

FAST AND ACCURATE PROPAGATION OF COHERENT LIGHT

RYAN D. LEWIS, GREGORY BEYLKIN, AND LUCAS MONZÓN

Abstract. We describe a fast algorithm to propagate, for any user-specified accuracy, a time-harmonic electromagnetic field between two parallel planes separated by a linear, isotropic, and homogeneous medium. The analytic formulation of this problem (circa 1897) requires the evaluation of the so-called Rayleigh-Sommerfeld integral. If the distance between the planes is small, this integral can be accurately evaluated in the Fourier domain; if the distance is very large, it can be accurately approximated by asymptotic methods. In the large intermediate region of practical interest

$$u(\mathbf{x}, z) = \frac{1}{z} \int_{\mathbb{R}^2} f(\mathbf{y}) \frac{e^{i2\pi R}}{R} d\mathbf{y}, \quad z > 0,$$

where $R = \sqrt{z^2 + |\mathbf{x} - \mathbf{y}|^2}$. Given the $u(\mathbf{y}, z)$ in the plane $z = 0$, we seek the $u(\mathbf{x}, z)$ for $z > 0$ that satisfies the free radiation condition. Expression (1) stands in contrast with the wave equation propagation

numerical simulation for propagation of a coherent wave packet in a dispersive medium. The numerical method is based on the split-step Fourier method [1].

In this paper, we propose a new numerical method for the propagation of a coherent wave packet in a dispersive medium. The method is based on the split-step Fourier method [1] and the fast Fourier transform (FFT) [2]. The method is simple and efficient, and it can be applied to a wide range of problems. The method is based on the split-step Fourier method [1] and the fast Fourier transform (FFT) [2]. The method is simple and efficient, and it can be applied to a wide range of problems.

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Preliminaries

The Rayleigh-Sommerfeld Formula. The behavior of a wave packet in a dispersive medium is governed by the Rayleigh-Sommerfeld formula [3]. The formula is given by

Given the boundary data $u(x, z=0)$ and $f(x)$ we rewrite (1) as

$$u(x, z) = \int_{\mathbb{R}^2} f(y) K_z(x-y) dy,$$

where the kernel $K_z(r)$ is given by

$$K_z(r) = \frac{e^{i2z\sqrt{1+(r/z)^2}}}{iz} - \frac{1}{r/z} - \frac{i}{z\sqrt{1+(r/z)^2}}, \quad r \geq 0.$$

Denote the Fourier transform of the boundary data as

$$f(p) = \int_{\mathbb{R}^2} f(x) e^{-i2x \cdot p} dx,$$

we write (1) in the Fourier domain as

$$u(x, z) = \int_{\mathbb{R}^2} f(p) K_z(p) e^{i2x \cdot p} dp,$$

where the Fourier transform of the kernel $K_z(r)$ is an entire function satisfying

$$K_z(p) = e^{i2z\sqrt{1-p^2}}, \quad |p| \leq 1/z.$$

Our goal is to validate a numerical scheme that the computational cost does not increase with the span z . It is clear that the spatial kernel $K_z(r)$ is a high oscillatory function of r when z is small and that the Fourier transform $K_z(p)$ is a high oscillatory function of p when z is large. For an physical interpretation, note that the span z in the interior at r on $K_z(r)$ and $K_z(p)$ are both high oscillatory and the numerical approximation of u using the Fourier transform is problematic. In § 2 we will show how to approximate the interior error and the numerical cost of the approximation. The numerical approximation of the boundary data $u(x, z=0)$ and $f(x)$ are assumed to be smooth propagation problems or interior at an arbitrary value of z . For small values of z it is well known that the problem is a boundary value problem and for large values of z the problem is a boundary value problem. § 2.1. A and B. Both on the support.

Remark 1. Given the non-regularity of the boundary data

$$u(x, z) = \int_{z=0} g(x),$$

where g is smooth or the unknown problem is

$$u(x, z) = \int_{\mathbb{R}^2} g(y) \frac{e^{i2z\sqrt{R^2 - |x-y|^2}}}{R} dy, \quad R = z^2 + |x-y|^2, \quad z > 0.$$

The numerical cost of our approach is also applicable to validation.

Slepian Functions. A phs ra st ust vntua a n spa an at th sa t ar ss nt a ban t nth Four r o an An appr at at a s rpt on o su h is was nt at b p an an h s o labor ators n b ons rn a spa t nan ban t n ra op rator an us n ts n un t ons to nt a lass o un t ons that hav ontro on ntr at on n both th spa an th Four r o ans p an t a show that th s nt ra op rator o ut s with th r nt ra op rator o lass a at a ph s s r b n th pro at sp h ro a wav un t ons both op rators shar th sa n un t ons

For our purposes we use n un t ons with ontro on ntr at on n a squar nth spat a o an an ban t to a s n th Four r o an h

The Unequally Spaced Fast Fourier Transform. In order to evaluate the forward transform, we consider

$$\sum_{m,m'=1}^M f_m y_{mm'} e^{i x \cdot y_{mm'}}$$

at output points $\mathbf{x}_{nn'} = \mathbf{x}_n, \mathbf{x}_{n'}$ where $n, n' = 1, \dots, N$. This is an N -point transform with an unequal spacing. A fast algorithm for this transform is given by the FFT algorithm with computational complexity $\mathcal{O}(N^2)$.

in \mathbf{z} is much larger than the spatial extent of $\mathbf{f}(\mathbf{y})$ so that a further approximation $e^{i\frac{\pi}{z}\|\mathbf{y}\|^2} \approx 1$ which when used in (1) leads to the

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$$K_z(r) = \frac{e^{i2z} e^{i\frac{\pi}{z}r^2}}{iz} A_z(r),$$

wh r

$$A_z(r) = \frac{1}{\sqrt{r/z}} \frac{i}{z \sqrt{r/z}} e^{i2z \left(\sqrt{1+(r/z)^2} - 1 - \frac{1}{2}(r/z)^2 \right)}.$$

where \mathbf{c} is the band vector to the input unit on \mathbf{f} in the band vector \mathbf{c}' with respect to the input rays or a similar expression using the quadrature form or

Let $\mathbf{y}_{mm'} = \mathbf{y}_m \cdot \mathbf{y}_{m'}$, $\mathbf{A}_{m,m'} = \dots, \mathbf{M}$ be the $\mathbf{M} \times \mathbf{M}$ tensor product of quadrature nodes with the corresponding quadrature weights $w_{m,m'}$ set an $\mathbf{N} \times \mathbf{N}$ matrix of output locations $\mathbf{x}_{nn'} = \mathbf{x}_n \cdot \mathbf{x}_{n'}$, $\mathbf{W}_{n,n'} = \dots, \mathbf{N}$ then apply the quadrature form or to the input rays and obtain an approximation to the output field at the source locations as

$$\mathbf{u}_{nn'} = \frac{e^{i2z} L}{iz} \sum_{m,m'=1}^M w_{m,m'} \mathbf{T}_{nn'mm'}^{(\cdot)} \cdot \mathbf{y}_{mm'} e^{i2 \ell \mathbf{x}_{nn'} \cdot \mathbf{y}_{mm'}}.$$

In the $\mathbf{N} \times \mathbf{N} \times \mathbf{M} \times \mathbf{M}$ fourth order tensors $\mathbf{T}^{(\cdot)}$

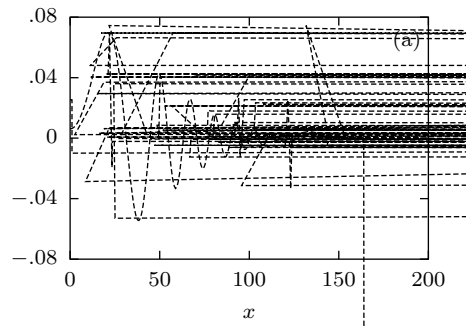
Lemma 4. Let $\mathbf{q}^{(l)}$, $\mathbf{U}_{nq}^{(l)}$, and $\mathbf{V}_{mq}^{(l)}$, where $l = 1, \dots, L$,

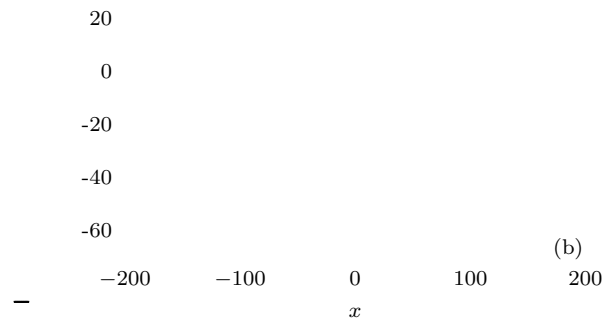
Theorem 5. The error of computing the field \mathbf{u} from $\mathbf{u}_{\text{nn}'}$ using $\mathbf{Q}_{\text{mm}'}$ is bounded by

$$|\mathbf{u}_{\text{nn}'}(z) - \mathbf{u}_{\text{nn}'}| \leq \frac{\mathbf{K} \mathbf{Q} \mathbf{R} \mathbf{f}_1}{z}.$$

The expression for $\mathbf{u}_{\text{nn}'}$ in (10) allows us to evaluate the error rapidly. First apply $\mathbf{Q}_{\text{mm}'}$ as a perturbation to the input signal $\mathbf{f}_{\text{mm}'}$.

to illustrate the relation between \mathbf{W}_{\max} and \mathbf{z}_{\min} or our theorem and \mathbf{W}'_{\max} and \mathbf{z}'_{\min} or the Fresnel approximation. It is obvious that $\mathbf{z}_{\min} \approx \mathbf{z}'_{\min}$ in the limit $\mathbf{z} \gg \lambda$.





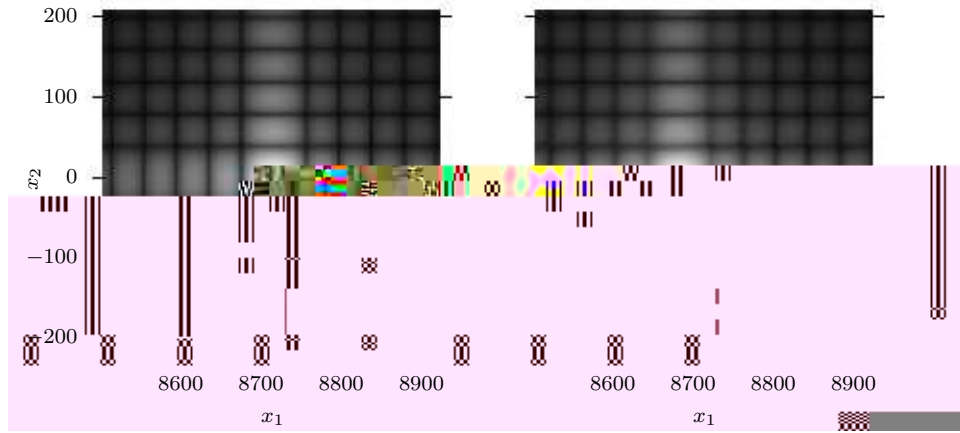
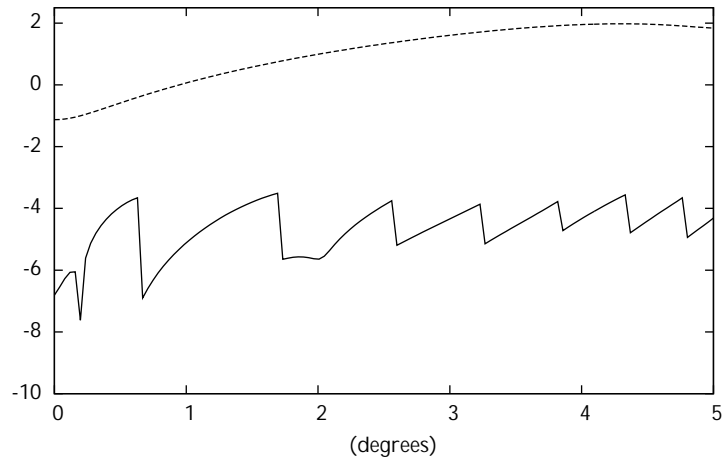


Figure 5.5. Comparison of the antenna of the Fourier transform of the optical axis. The top part shows the antenna of the Fourier transform of the optical axis. The bottom part shows the antenna of the Fourier transform of the optical axis. The top part shows the antenna of the Fourier transform of the optical axis. The bottom part shows the antenna of the Fourier transform of the optical axis.



in Figure 10, the trajectory of the spot is shown. In this case, the spot is not stationary. The trajectory of the spot is shown in Figure 10. The spot is shown to be moving in a circular path. The spot is shown to be moving in a circular path. The spot is shown to be moving in a circular path.

Representative Examples of Computational Cost. In the following, we present representative examples of computational cost.

Conclusions

have shown that the fast Fourier transform method is a powerful tool for the analysis of wave fields. In contrast to the traditional Fourier transform, the fast Fourier transform is a highly efficient algorithm for the analysis of wave fields. As a result, the fast Fourier transform is a powerful tool for the analysis of wave fields. The fast Fourier transform is a powerful tool for the analysis of wave fields. The fast Fourier transform is a powerful tool for the analysis of wave fields.

Acknowledgments

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